

# **The Nanoelectronic Modeling Tool (NEMO) and its Expansion to High Performance Parallel Computing**

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**Jet Propulsion Laboratory/California Institute of Technology**

**<sup>1</sup>University of Alabama in Huntsville**

**<sup>2</sup>University of Illinois**

# Outline

- **NEMO genealogy and general features**
- **Structural device optimization:**
  - **Massively parallel genetic algorithm package**
- **Full band transport simulations:**
  - **Electron transport**
  - **Hole transport**
- **Comparison between k.p and sp<sup>3</sup>s\***
- **Conclusion**

# NEMO Genealogy

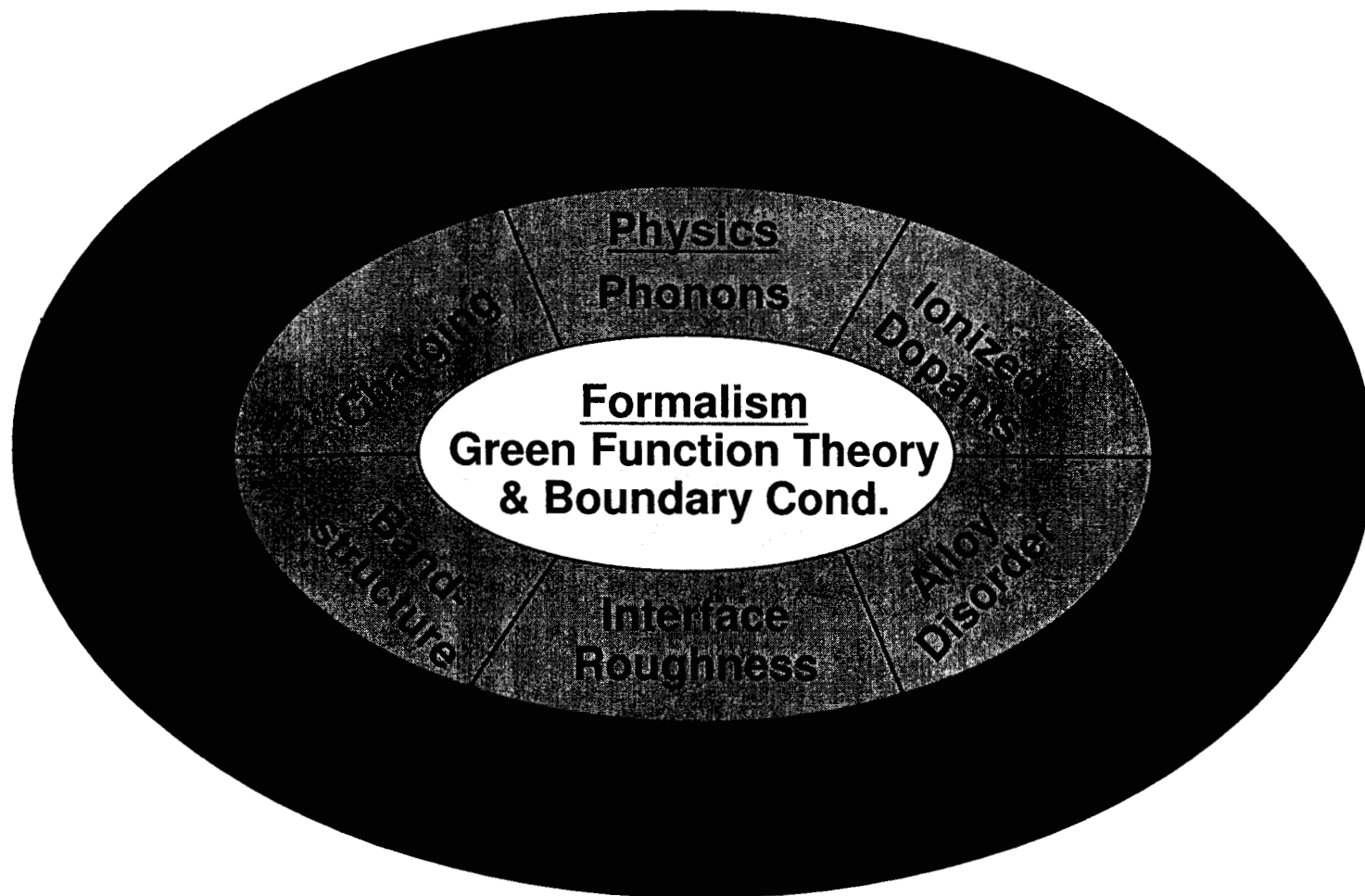
- **NEMO was developed under a government contract at Texas Instruments and Raytheon from 1993-1997**
  - **Theory**
    - **Roger Lake, Chris Bowen, Tim Boykin (UAH), GK**
  - **Graphical User Interface**
    - **Dan Blanks, GK**
  - **Programming Approach, Philosophy, and Prototypes**
    - **Bill Frensley (UTD), GK**
  - **Coding**
    - **Manhua Leng (UTD), Chenjing Fernando, Paul Sotirelis, Dejan Jovanovic, Mukund Swaminathan (UTA), GK**
  - **Experiments for verification**
    - **Ted Moise, Alan Seabaugh, Tom Broekaert, Berinder Brar, Yung-Chung Kao**
- **NEMO is based on non-equilibrium Green functions, in an implementation that is novel. The development of NEMO has benefited from the vast research on resonant tunneling diodes that had been done before the project.**

## Summary of NEMO Capabilities

- **Interface / Users:**
  - **FAST** and dirty design. **interactive**
  - **Comprehensive analysis (SLOW).** **batch**
- **Physics**
  - **Charging**
    - Semi-classical self-consistency, quantum self-consistency
  - **Scattering**
    - Phonons, alloy disorder and interface roughness (1band)
  - **Bandstructure**
    - 1, 2, 10 tight binding band models nearest and next nearest neighbor coupling.
  - **Realistically long devices**
    - Novel boundary conditions

**NEMO can trade of CPU time and memory against a variety of models.**

# All of NEMO's Facets: Formalism, Physics, and Technology



Approximately 250,000 lines of code at time of delivery

# **“Genetically Engineered” Nanoelectronic Devices**

## **Objective:**

- Automated device synthesis and analysis using genetic algorithms.
- Material spectroscopy through genetic algorithms analysis.

## **Justification:**

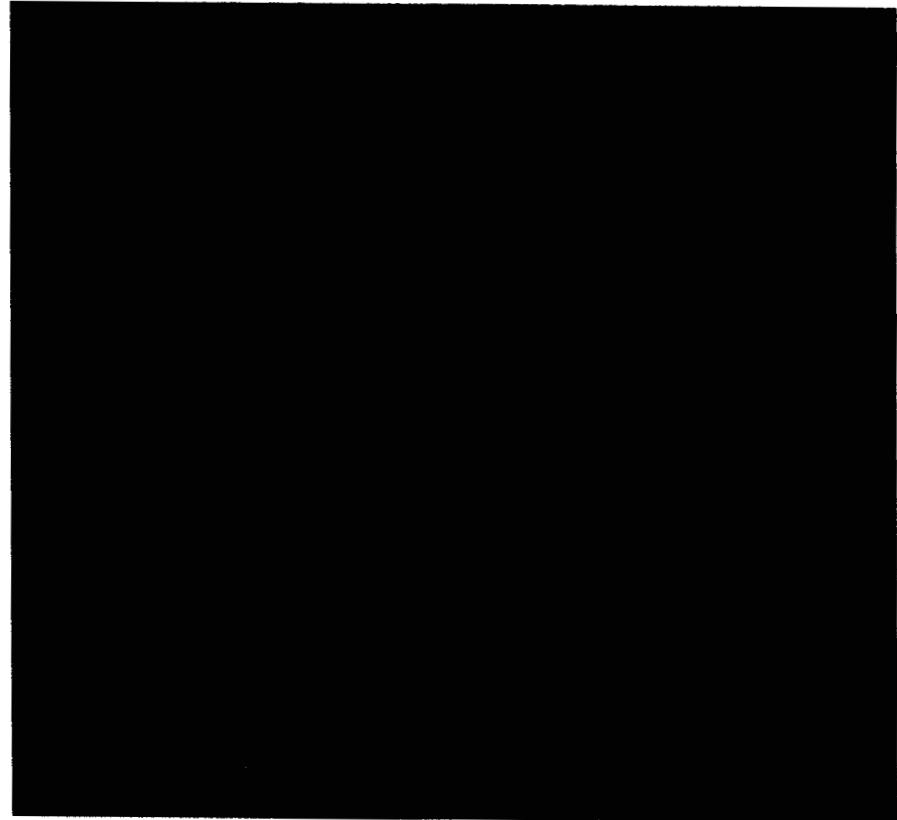
- Empirical Design (usual process) is sub-optimal. Complete design space search is unfeasible.  
=> Develop automated design tools.

## **Impact:**

- Rapid nanotechnology device synthesis and development.
- Generation of novel devices.

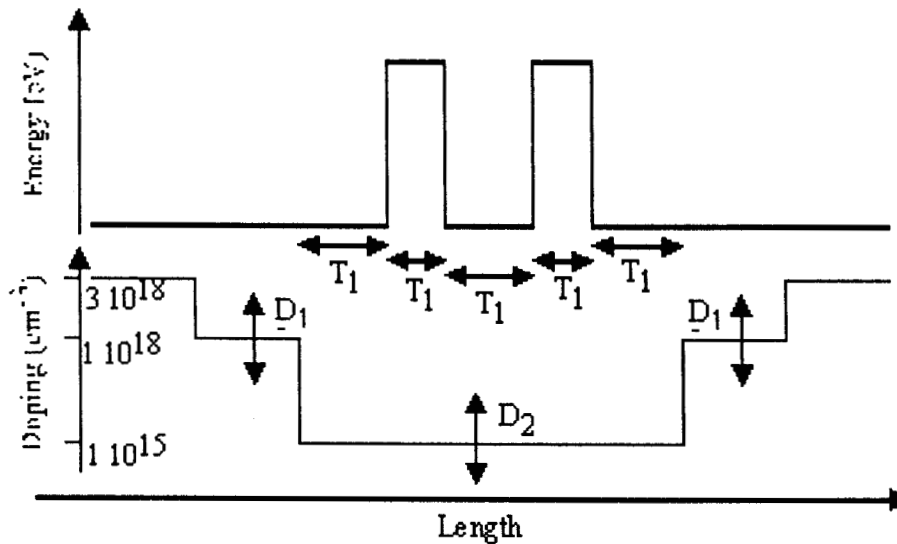
## **Approach:**

- Augment NEMO to analyze individual structures in parallel.
- Augment parallel genetic algorithm package (PGApack) to optimize and select desired structures in NEMO.



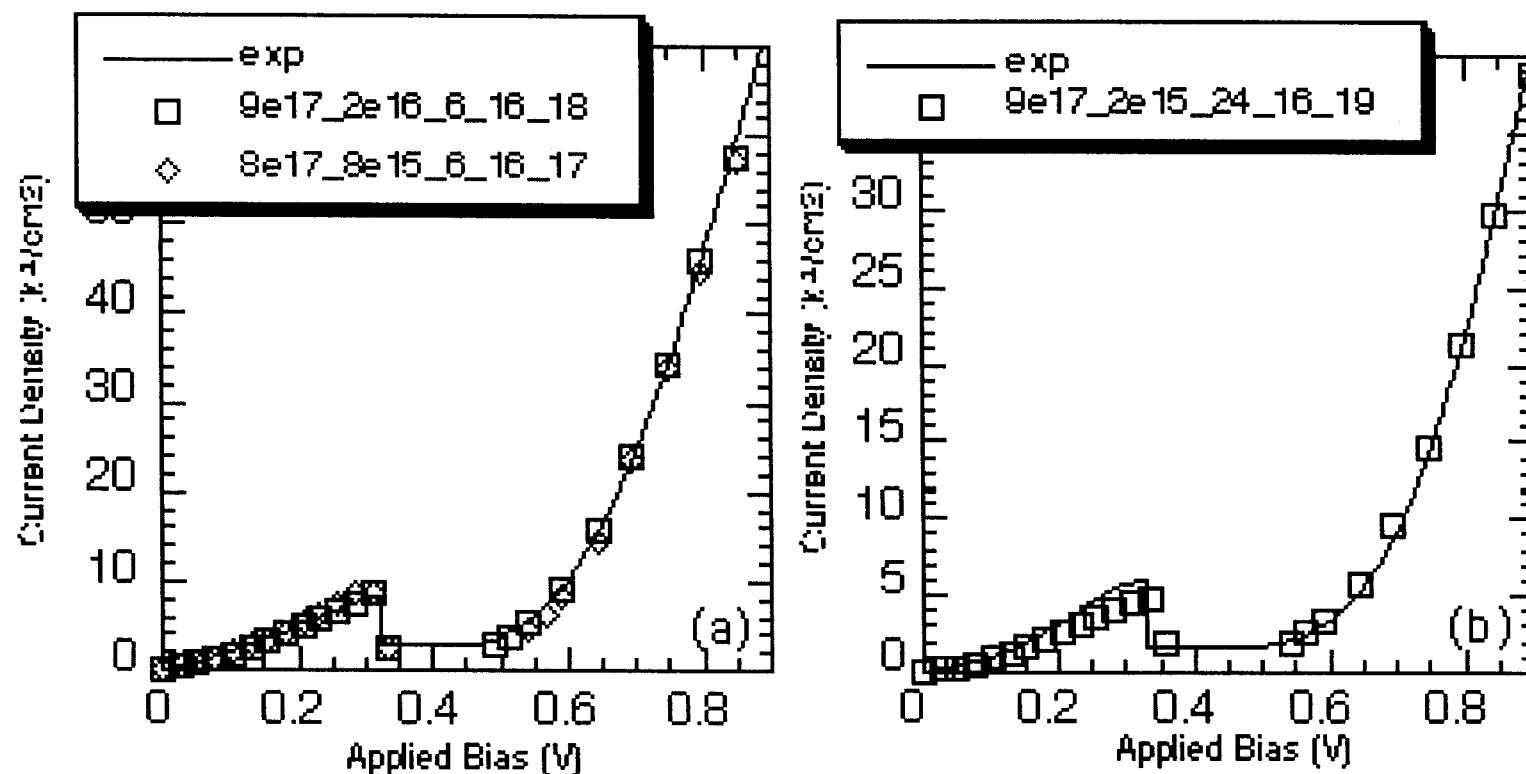
Proposed system architecture.  
Prototype is operable in batch mode.

# First Simulation Results: Structural Analysis



- Allow genetic algorithm to vary 5 different structural parameters:
  - 3 Thicknesses: well, barrier, spacer
  - 2 Dopings: low doped spacer, unintentional doping in center
- Employ parameterized non-parabolic single band model with full quantum charge self-consistency and transverse momentum integration.
- Developed fitness function for typical RTD I-V curves. Need to shoot for peak position and amplitude, slope at peak and relative and absolute errors.

# First Simulation Results: Structural Analysis

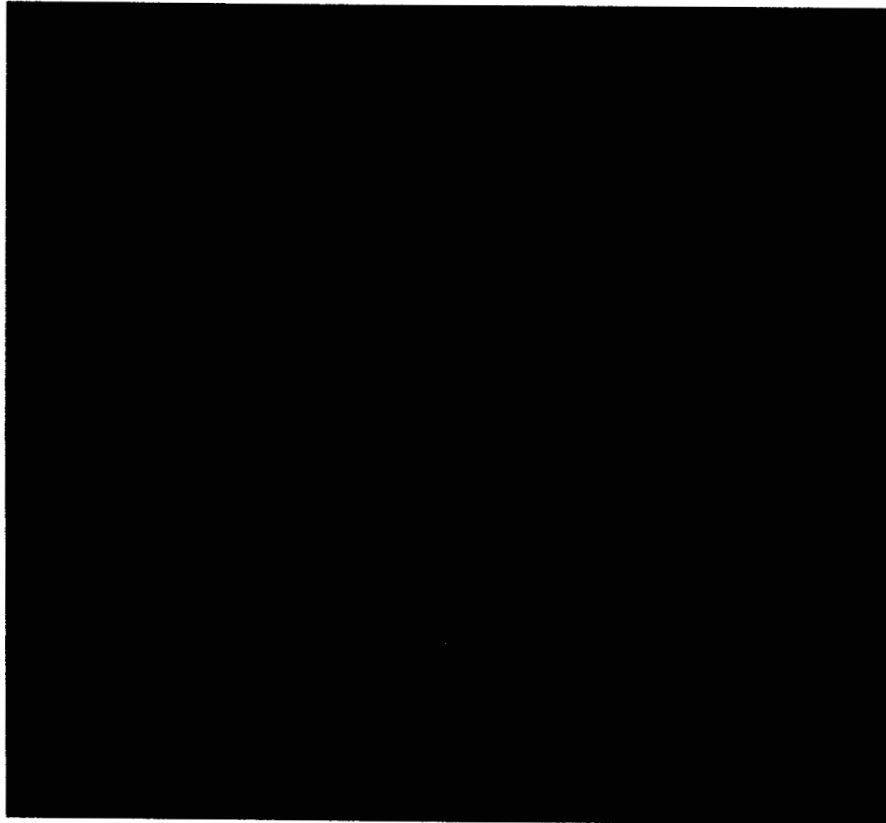


• Analyzed two similar InGaAs/InAlAs structures :

- RTD 1:  $D1=1e18/cm^3$ ,  $D2=1e15/cm^3$ ,  $T1=7ml$ ,  $T2=16ml$ ,  $T3=16ml$
- RTD 2:  $D1=1e18/cm^3$ ,  $D2=1e15/cm^3$ ,  $T1=20ml$ ,  $T2=16ml$ ,  $T3=16ml$



# Future Interest



- Analyze material parameter influence on overall device performance  
-> material spectroscopy
- Implement general architecture such that a variety of different simulation tools can be plugged into an optimization tool.
- Explore other optimization algorithms, such as simulated annealing or directive approaches within the same framework.  
  
-> scripting tools that can link different tools  
  
-> Tcl/Tk

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## Transport via Transmission Coefficients

$$I \propto \int dk_x \int dk_y \int dE T(E, k_x, k_y) (f_L(E) - f_R(E))$$



**Cylindrical Coordinates**

$$I \propto \int d\varphi \int k dk \int dE T(E, k, \varphi) (f_L(E) - f_R(E))$$



**Throw out angular dependence**

$$I \propto 2\pi \int k dk \int dE T(E, k) (f_L(E) - f_R(E))$$

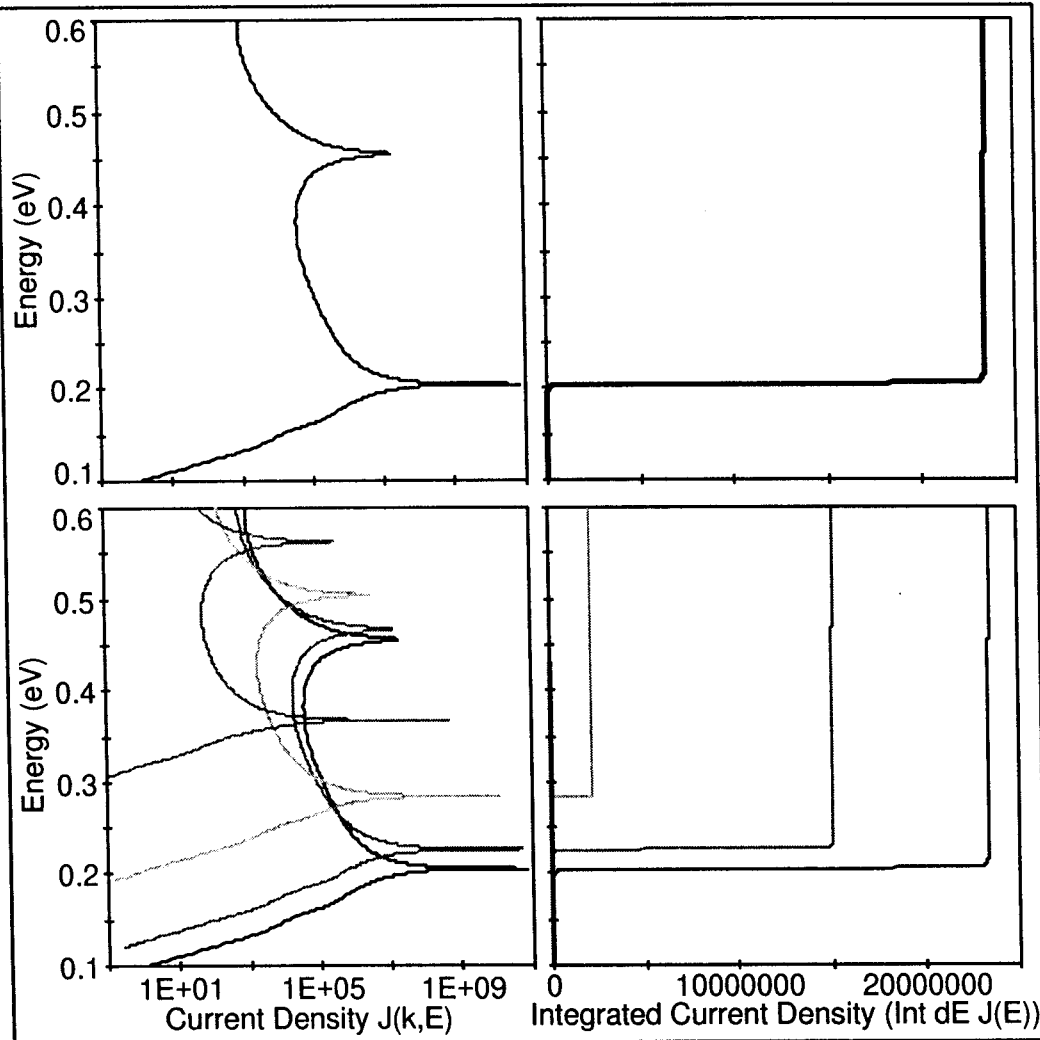


**Parabolic transverse subbands**

$$I \propto \rho_{2D} \int T(E) (f_L(E) - f_R(E))$$

## **Full Band Simulation of RTD Electron Transport**

- **Mechanics of an energy  $E$  and transverse momentum  $k$  integration.**
- **Unphysical effects when parabolic subbands are assumed:**
  - **Overshoot at RTD turnoff.**
  - **Spurious bistability in charge self-consistent simulation.**
- **Transverse subbands.**
- **Full band simulation eliminates these spurious simulation effects.**



$$J(E, k)$$

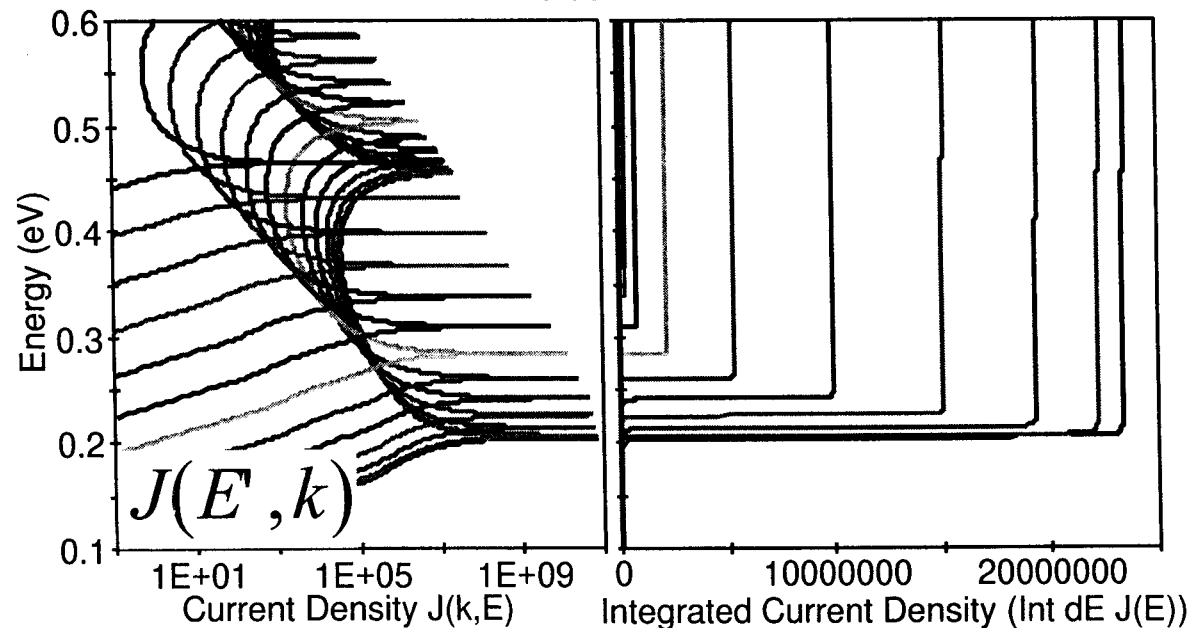
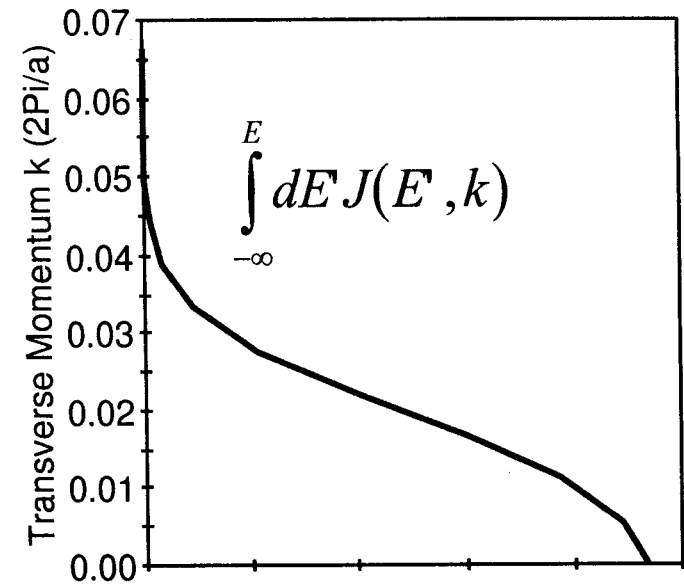
$$\int_{-\infty}^E dE J(E, k)$$

## Mechanics of 2D Integration

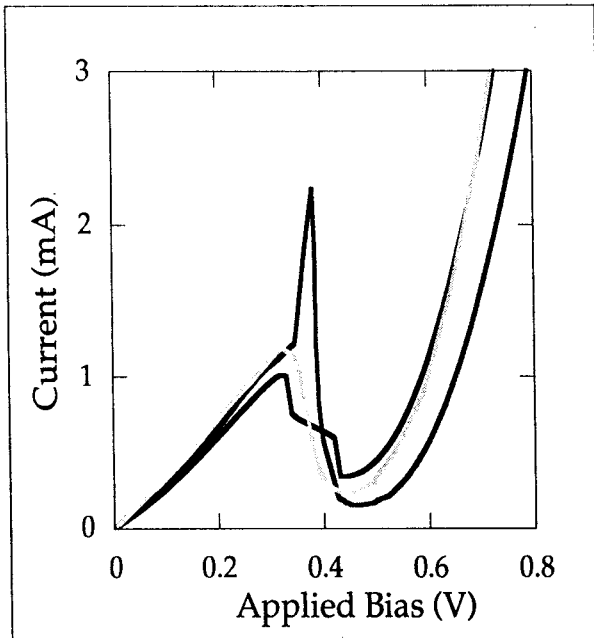
- Transmission coefficient is masked by Fermi distribution in injecting lead.
- Running sum integral points out where in energy space significant current contributions occur.
- Multiple instead of a single transmission coefficient are evaluated and summed up.

# Mechanics of 2D Integration

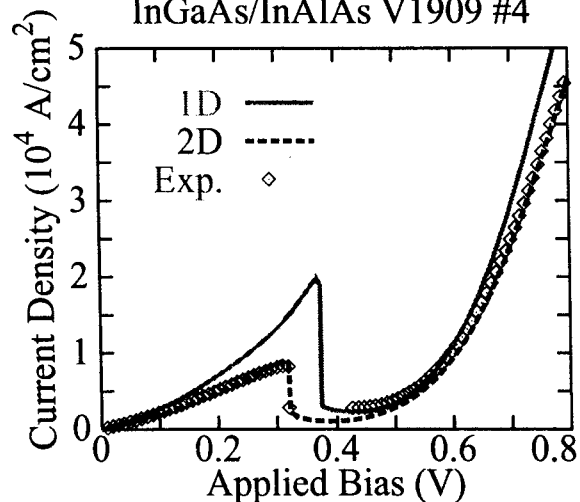
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# Full Band Simulation of Electron Transport



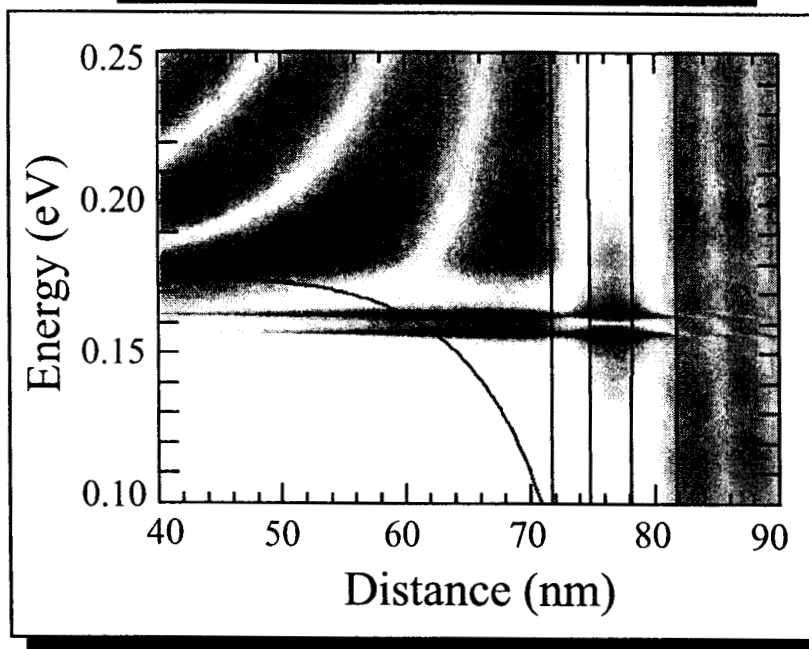
InGaAs/InAlAs V1909 #4



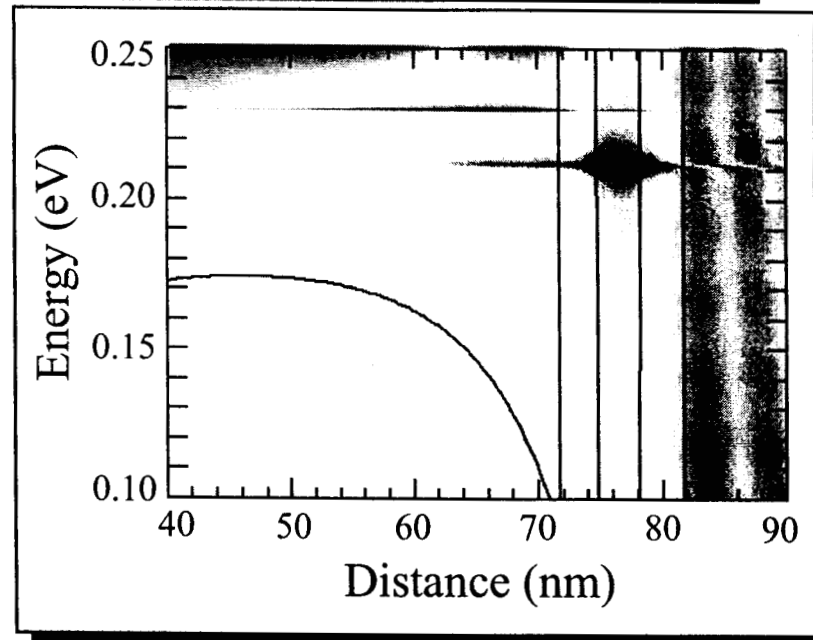
- 1D integration assuming parabolic subbands can lead to unphysical current overshoots.
- 2 Examples on InGaAs/InAlAs simulations:
  - Sp3s\* simulation with partial charge self-consistency  
-> sharp spike at turn-off
  - Parameterized single band simulation which incorporates the band-non-parabolicity  
-> overall current overshoot.
- -> 2D integration fixes these unphysical results.

# Resonance Coupling vs. Transverse Momentum

Density of States ( $k_x=0.00$ )



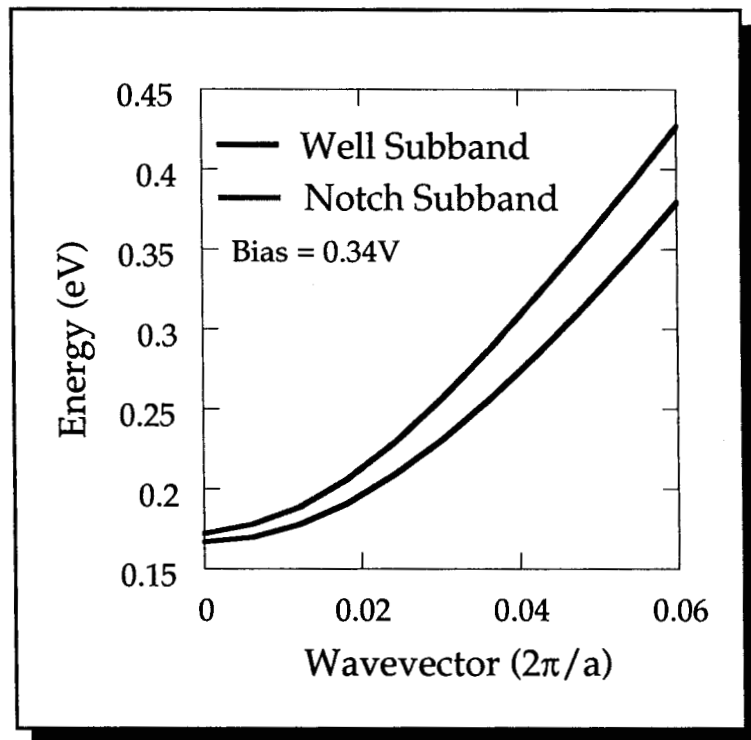
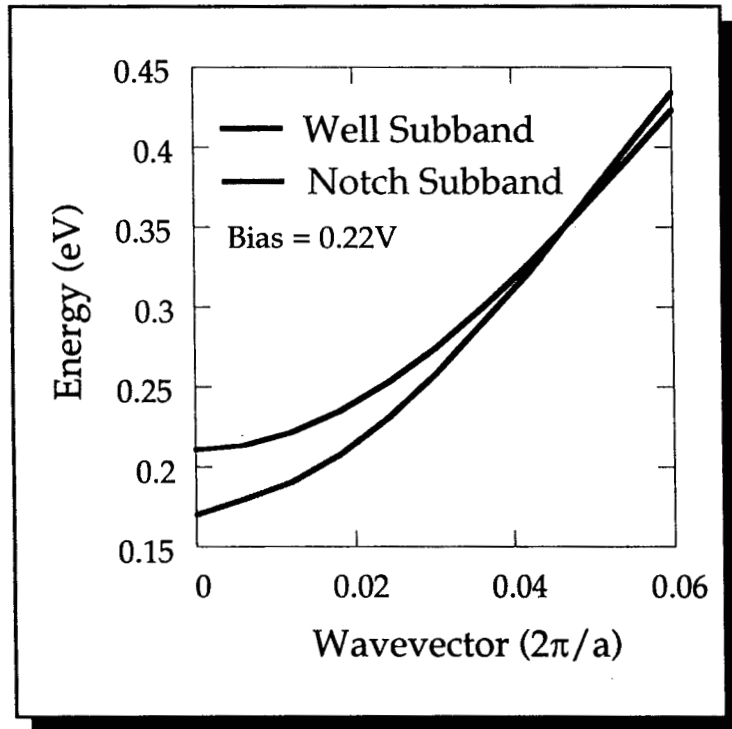
Density of States ( $k_x=0.03$ )



**Resonance coupling depends on the transverse momentum**

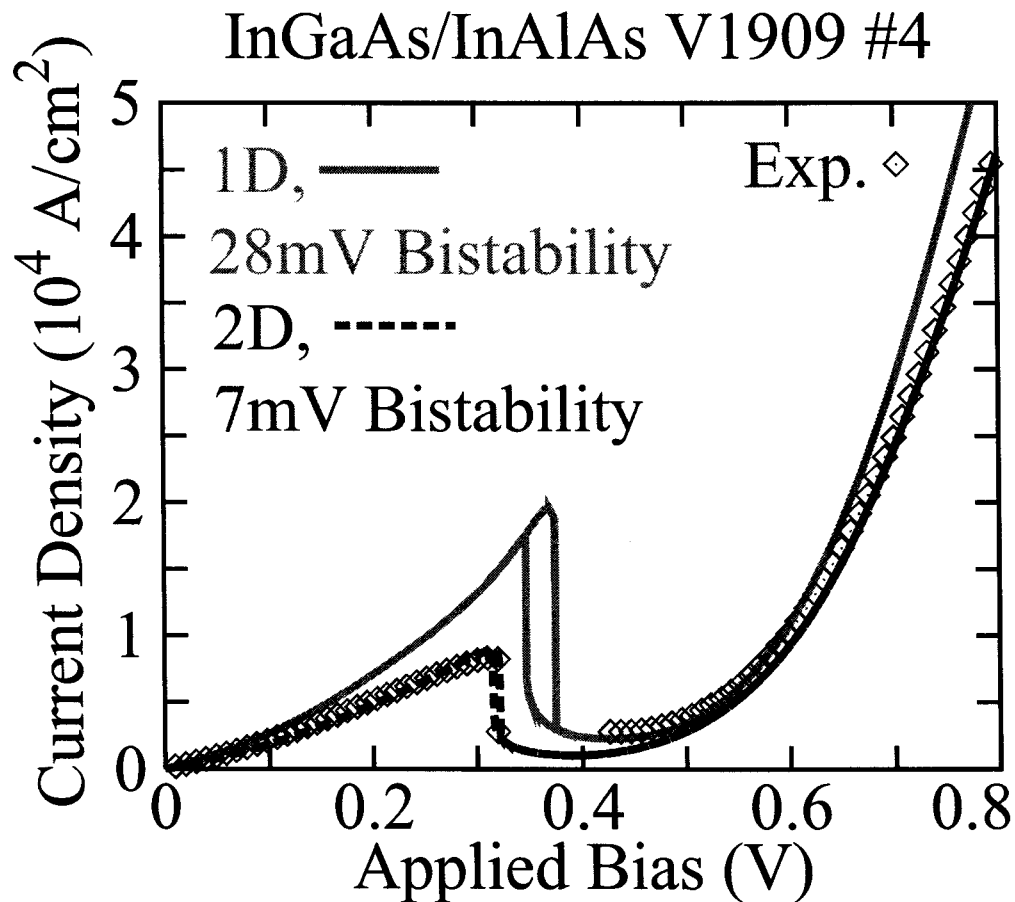


# Quantum Well and Notch Subbands



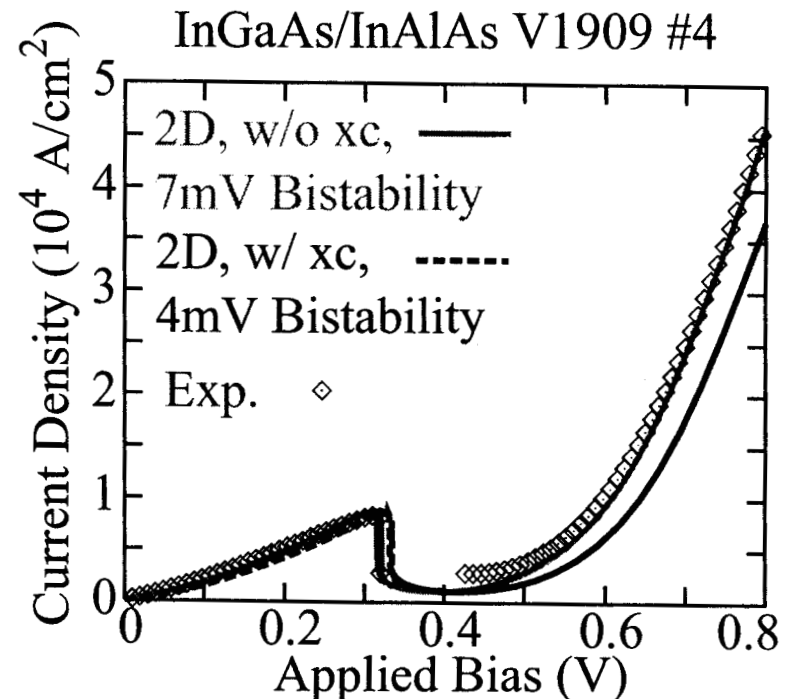
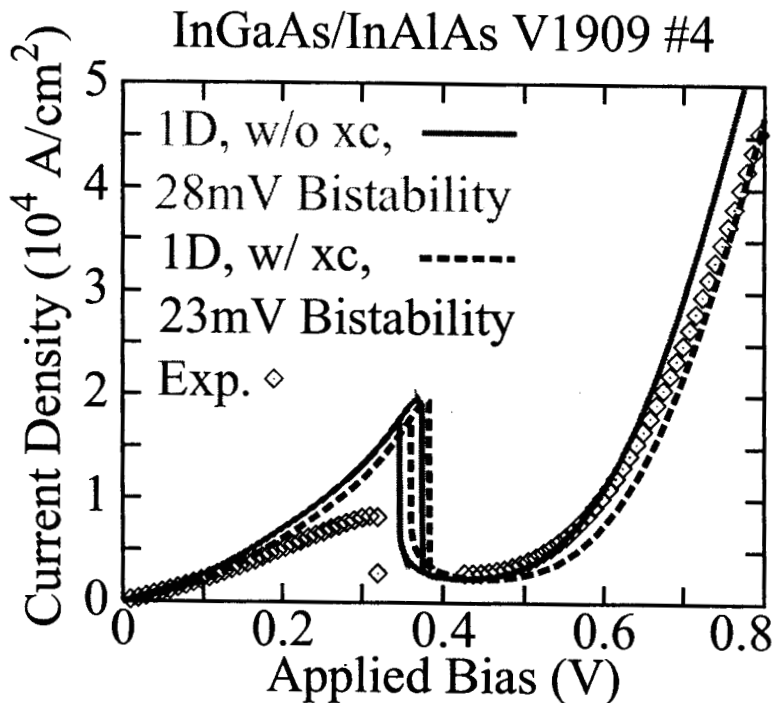
- The dispersions are non-parabolic
- There is no “perfect” overlap of the subbands

# Spurious Bistability



- Most quantum charge self-consistent simulations of RTD's exhibit a bistability in the NDR region when the simulation is performed in a reverse voltage sweep.
- This is (we believe) a numerical artifact and typically not observed in experiments.
- Full band integration reduces the spurious bistability significantly.

# Effect of Exchange and Correlation Potential



- Calculate the exchange and correlation potential in the local density approximation.
- Exchange and correlation energy does not eliminate (in general) the bistability, it does reduce it however.
- Inclusion of scattering in the simulation reduces the bistability region as well.

# Outline

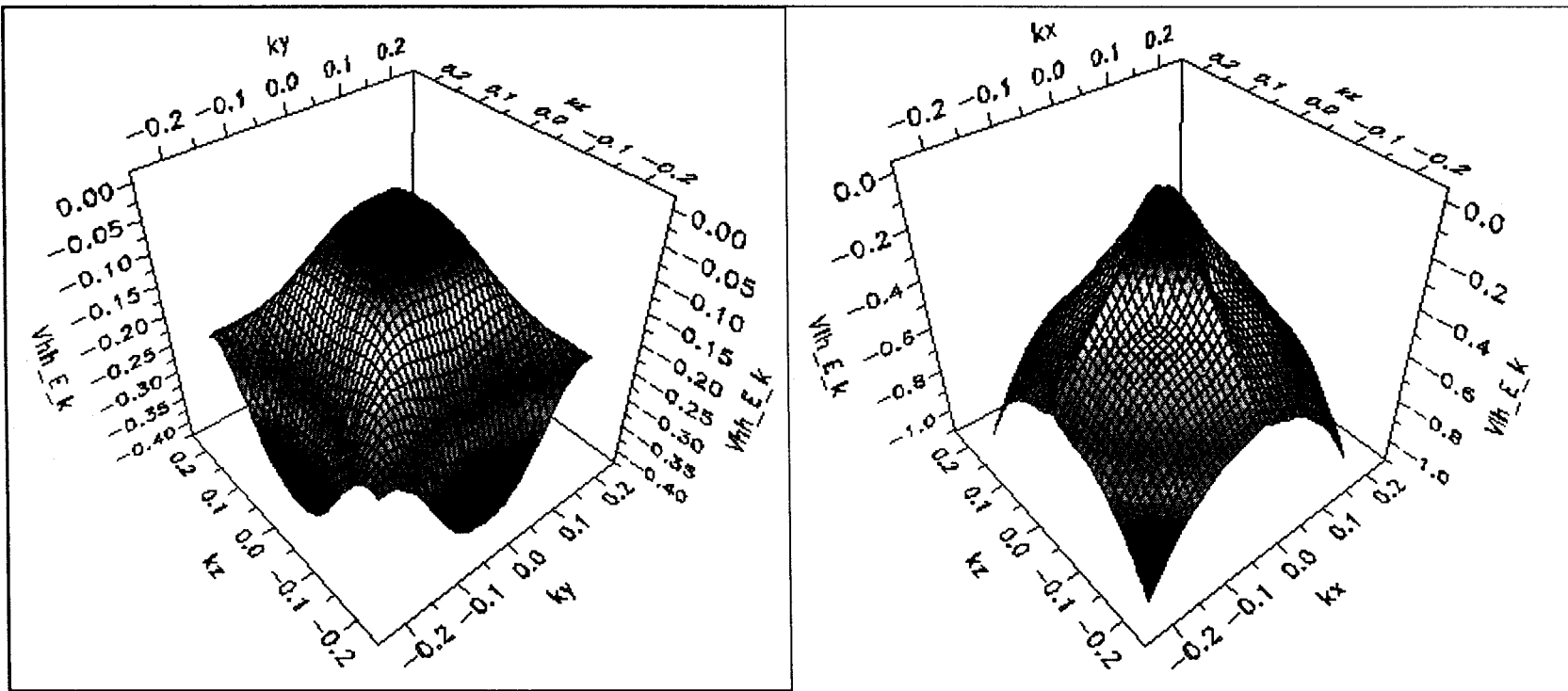
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## **Hole Transport in a AlAs/GaAs/AlAs RTD**

- **Resonance states**
- **Transverse momentum subbands**
- **Need for energy and transverse momentum integration**
- **Current flow away from zone center**
- **Dependence on the direction of the transverse momentum ( $\langle 100 \rangle$  vs.  $\langle 110 \rangle$  integration).**

# GaAs Bulk HH and LH Bandstructure

Computed in second nearest neighbor sp<sup>3</sup>s\* tight binding model with explicit spin (20x20 basis).



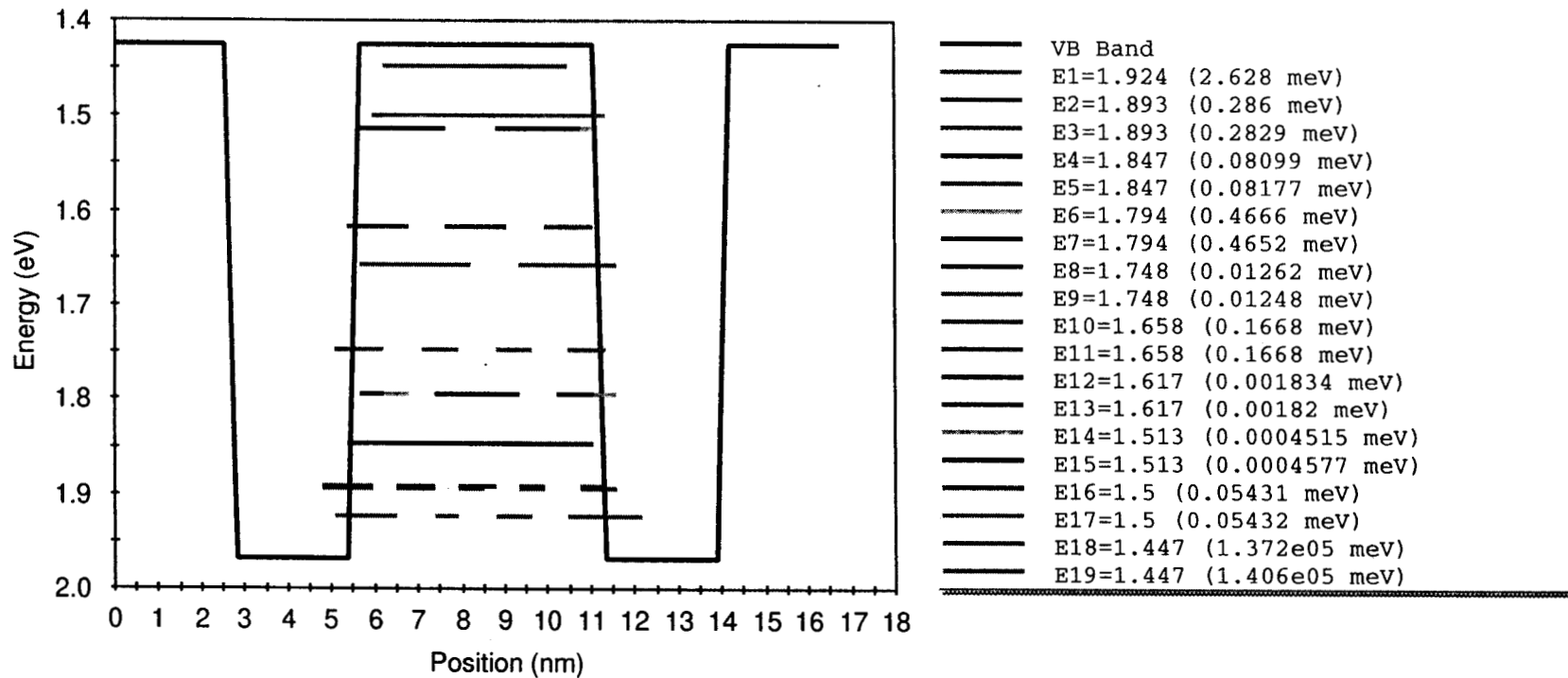
Heavy holes heavier in  $\langle 110 \rangle$  than  $\langle 100 \rangle$

$$m_{110} > m_{100}$$

Light holes lighter in  $\langle 110 \rangle$  than  $\langle 100 \rangle$

$$m_{110} < m_{100}$$

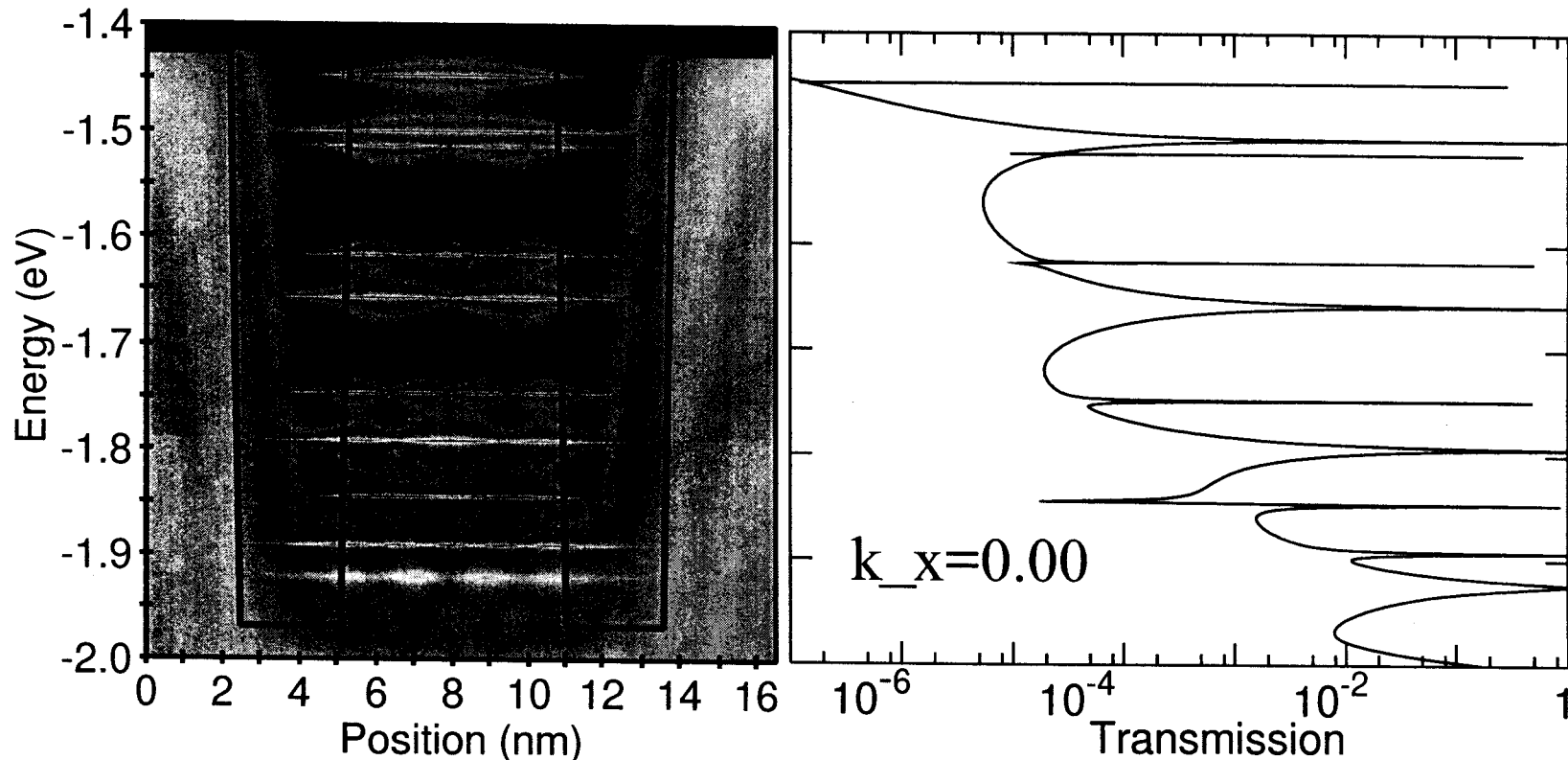
# Hole Resonance States in GaAs/AlAs RTD



- Look at simple structure:
  - 10 monolayer AlAs barriers
  - 20 monolayer GaAs well
  - Flat band conditions

- Use 2nd nearest neighbor  $sp^3s^*$  tight binding model
- Compute resonance energies and resonance linewidths using a order N non-hermitian matrix eigenvalue solver

# Hole Transport in a GaAs/AlAs RTD



- **Density of states (DOS):**

- Low DOS - Dark (blue) tones
- High DOS - light (red) tones

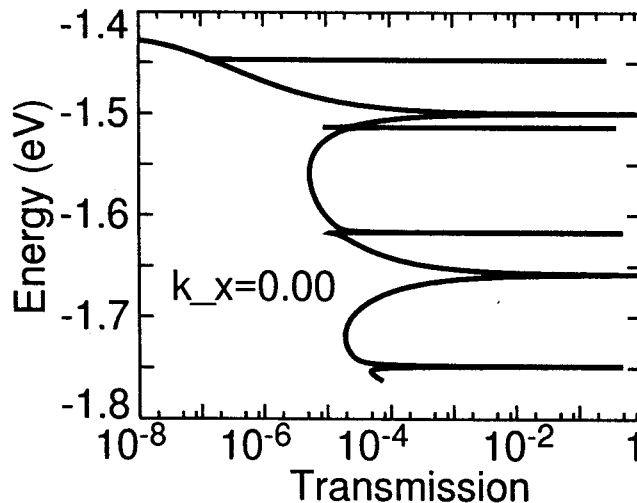
- **Transmission coefficient at  $k_x=0$**

- Light hole states strongly coupled to continuum  $\rightarrow$  wide resonances
- Heavy hole states weakly coupled to continuum  $\rightarrow$  narrow

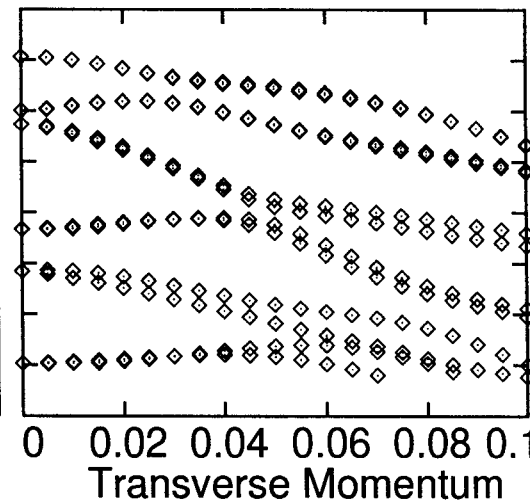


# Transverse Hole Subbands in GaAs/AlAs RTD

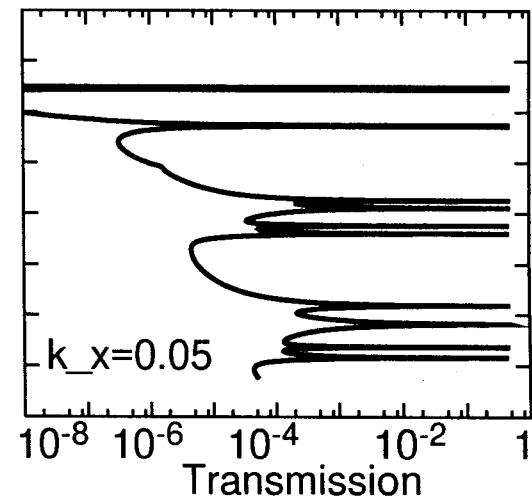
- Transmission coefficient at  $k_x=0$



- Resonance states as a function of transverse momentum:  
→ transverse subbands

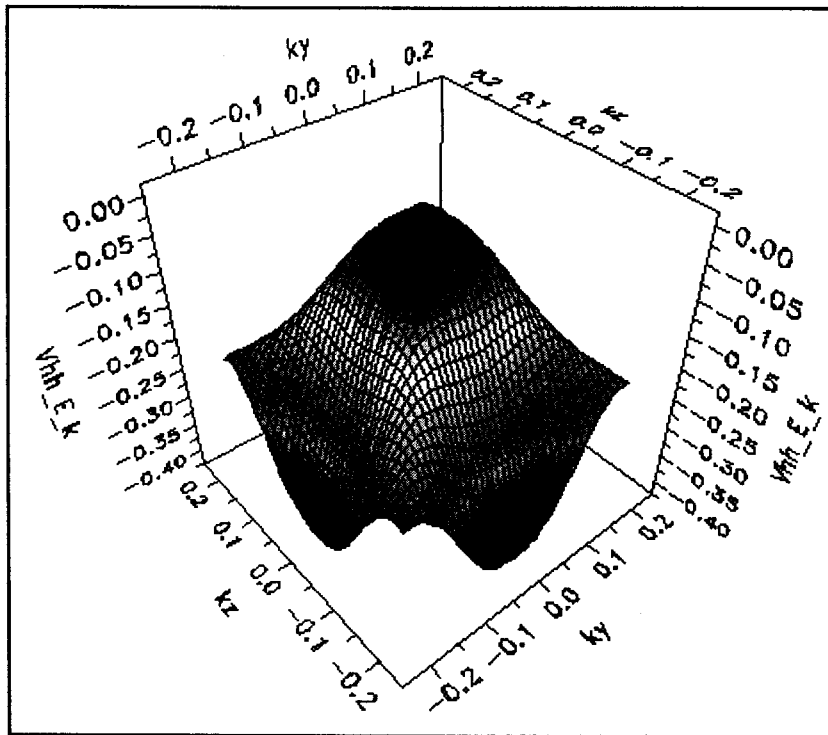


- Transmission coefficient at  $k_x=0.05$

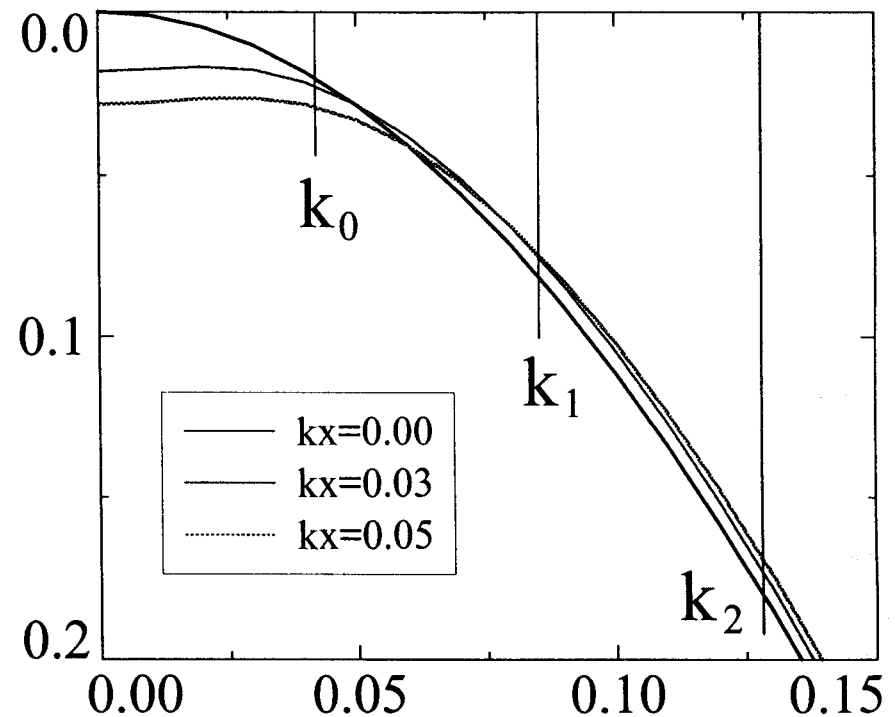


- Transverse subbands exhibit:
  - complex structure of anti-crossings
  - Non-monotonic behavior - some subbands are electron-like.
- Why do some of the subbands increase in energy?

# Excited HH-States Move up in Energy

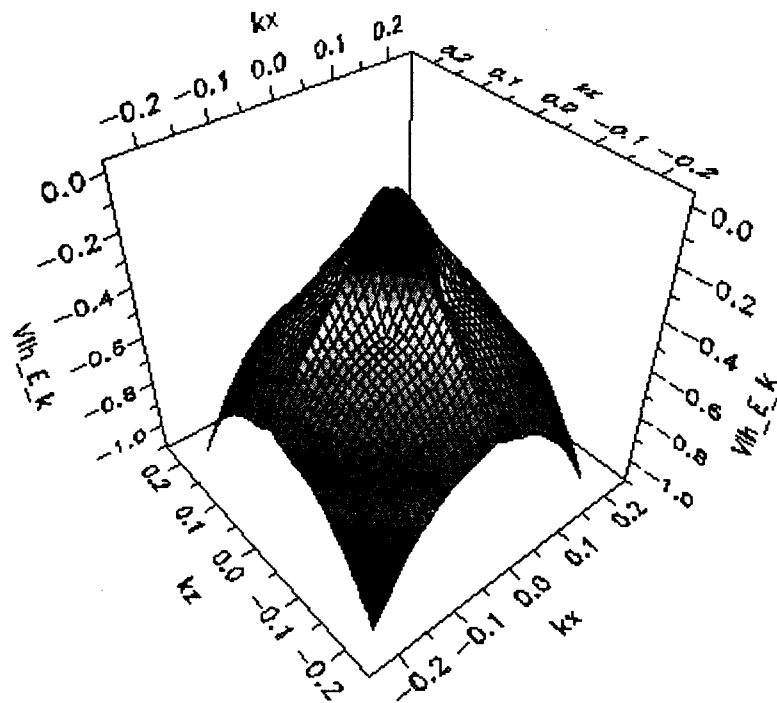


- HH Bulk dispersion in  $(k_x, k_y, 0.0)$  shows strong anisotropy in  $\langle 100 \rangle$  vs.  $\langle 110 \rangle$  direction.

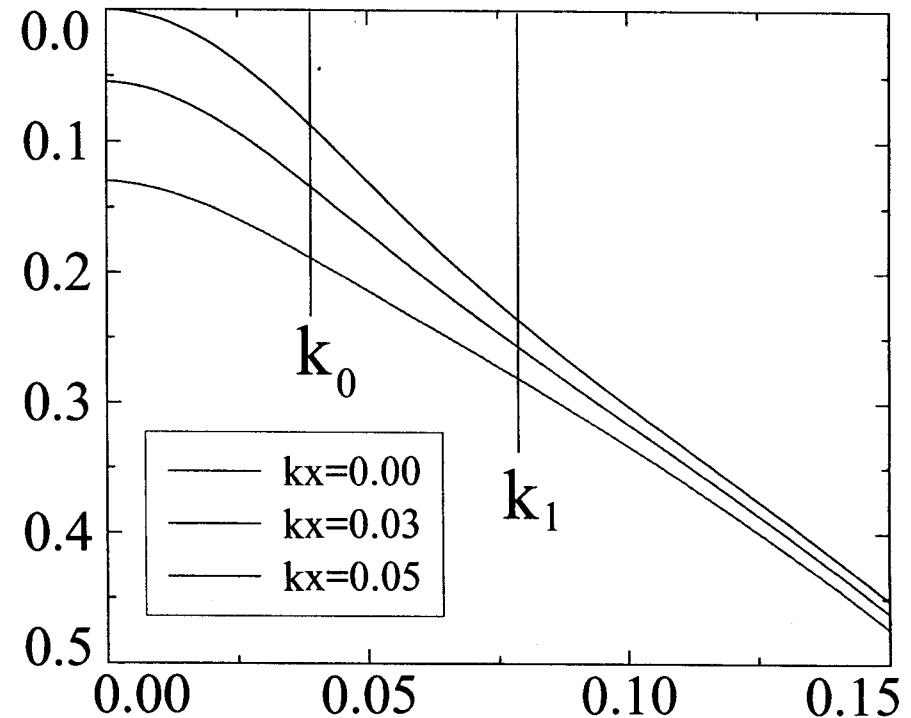


- Slice  $(k_x, k_y, 0.0)$  surface for 3 different  $k_x$  (black, blue, and red curve).
- $k_0$  indicates the heavy hole ground state determined by  $\sim \pi/L$
- $k_1 = 2k_0$  : first excited state - moves up
- $k_2 = 3k_0$  : second excited state - moves up

# LH-States Move down in Energy

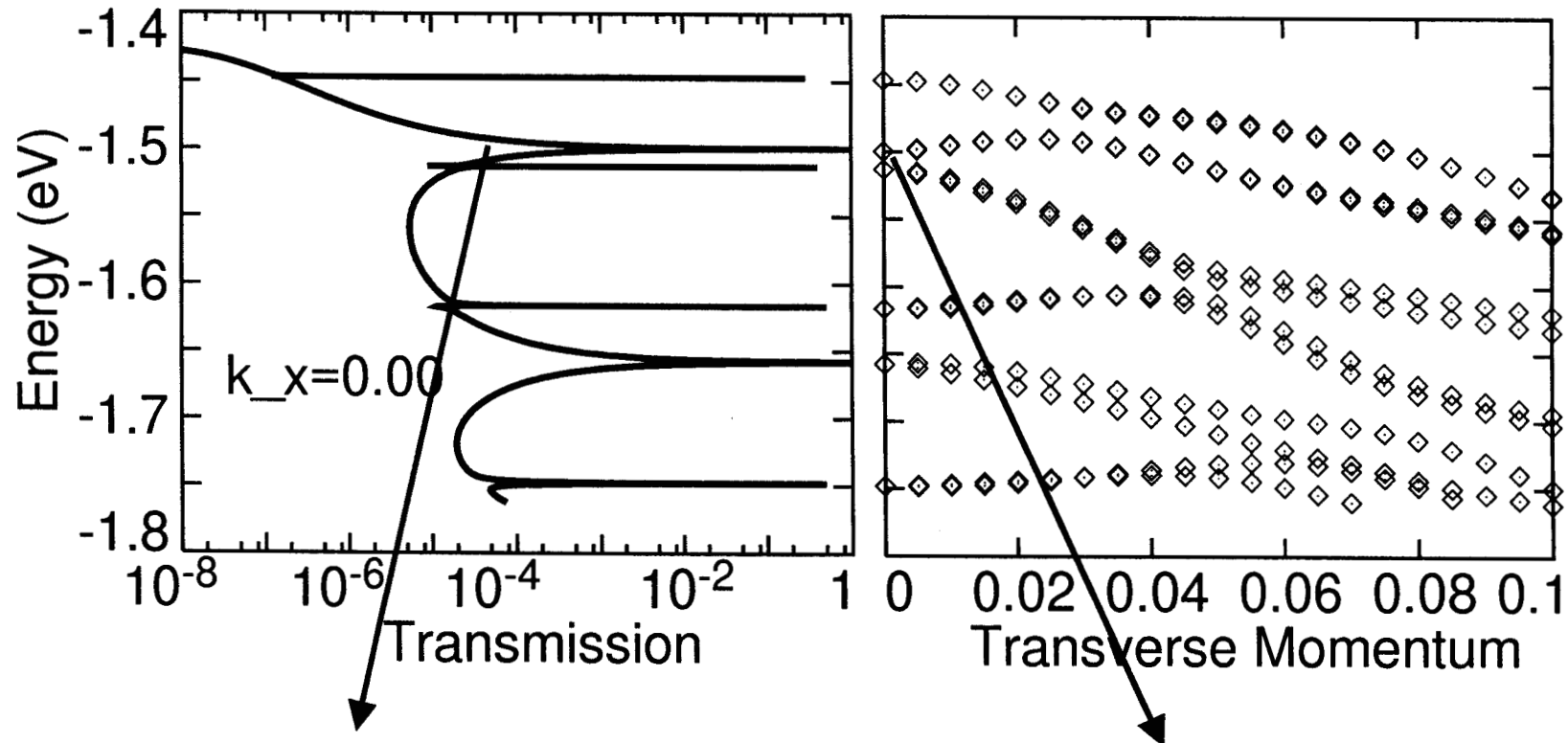


- HH Bulk dispersion in  $(k_x, k_y, 0.0)$  shows strong anisotropy in  $\langle 100 \rangle$  vs.  $\langle 110 \rangle$  direction.



- Slice  $(k_x, k_y, 0.0)$  surface for 3 different  $k_x$  (black, blue, and red curve).
- $k_0$  indicates the light hole ground state determined by  $\sim \pi/L$   
ground state moves down fast
- $k_1 = 2k_0$ : first excited state moves down

# Cave Canem! Beware of the Wolf!

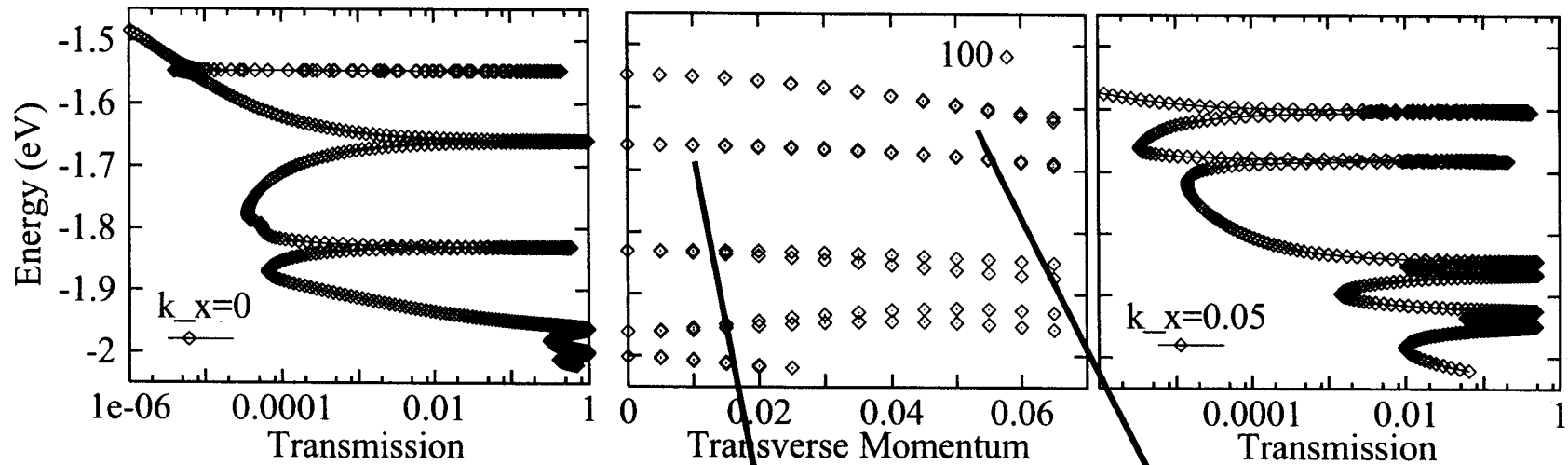


- What looks like a light hole resonance with a neighboring heavy hole resonance is already a strongly mixed state

- Observe an anticrossing of two states at  $k_x = 0$

More Latin: nemo = nobody

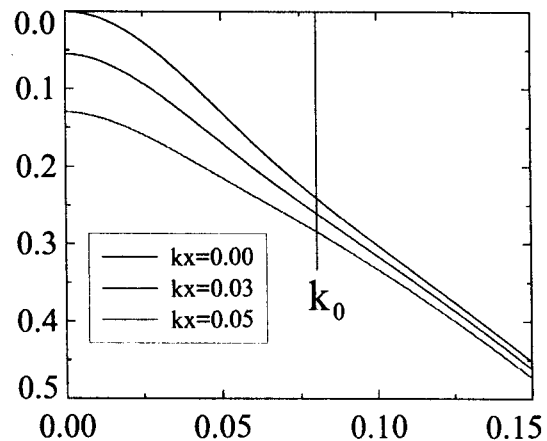
# Narrow GaAs Well



- AlAs/GaAs/AlAs structure, 10/07/10 ml thickness
- LH1 at  $k_x=0$  is de-coupled from HH2 due to strong confinement

LH dispersion is almost flat

LH1 at  $k_x=0.05$  anti-crosses with HH1



## Hole Transport

$$I \propto \int dk_x \int dk_y \int dE T(E, k_x, k_y) (f_L(E) - f_R(E))$$



**Cylindrical Coordinates**

$$I \propto \int d\varphi \int k dk \int dE T(E, k, \varphi) (f_L(E) - f_R(E))$$



**Throw out angular dependence**

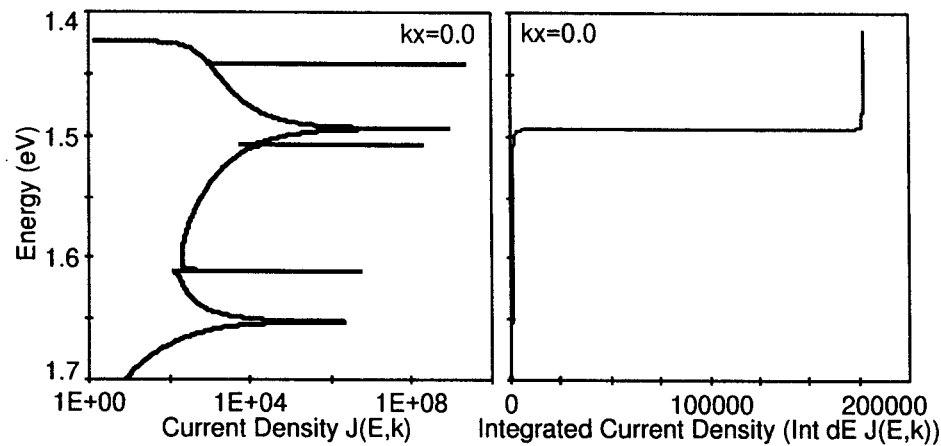
$$I \propto 2\pi \int k dk \int dE T(E, k) (f_L(E) - f_R(E))$$



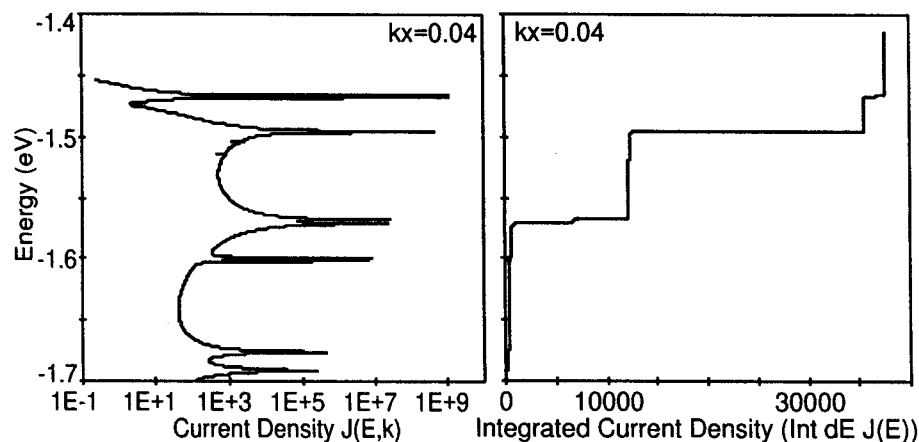
**Parabolic transverse subbands**

$$I \propto \rho_{2D} \int T(E) (f_L(E) - f_R(E))$$

# Current Integral Varies Qualitatively with Different Transverse Momenta



**$k=0.00$ :**  
current flows at one energy

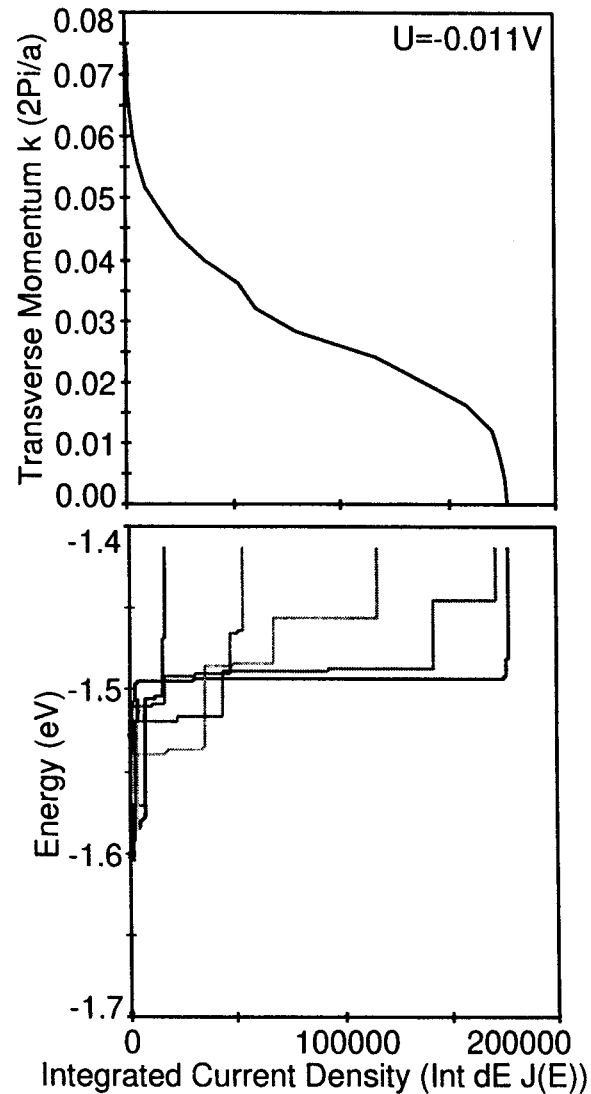


**$k=0.04$ :**  
current flows at multiple energies

$$J(E, k)$$

$$\int_{-\infty}^E dE J(E, k)$$

# Current Density $J(k)$

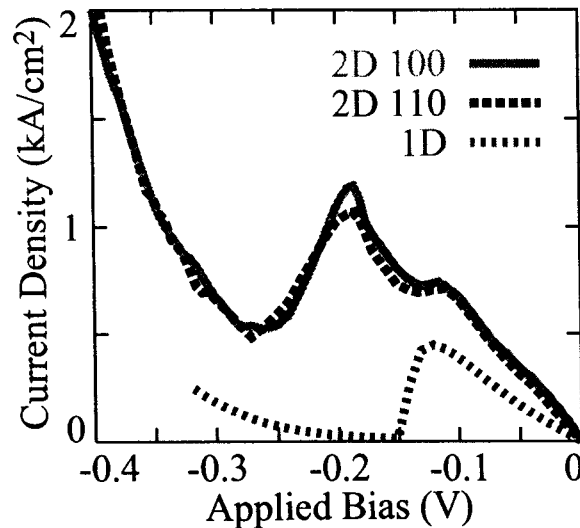


$$J(k) = \int_{-\infty}^{\infty} dE J(E, k)$$

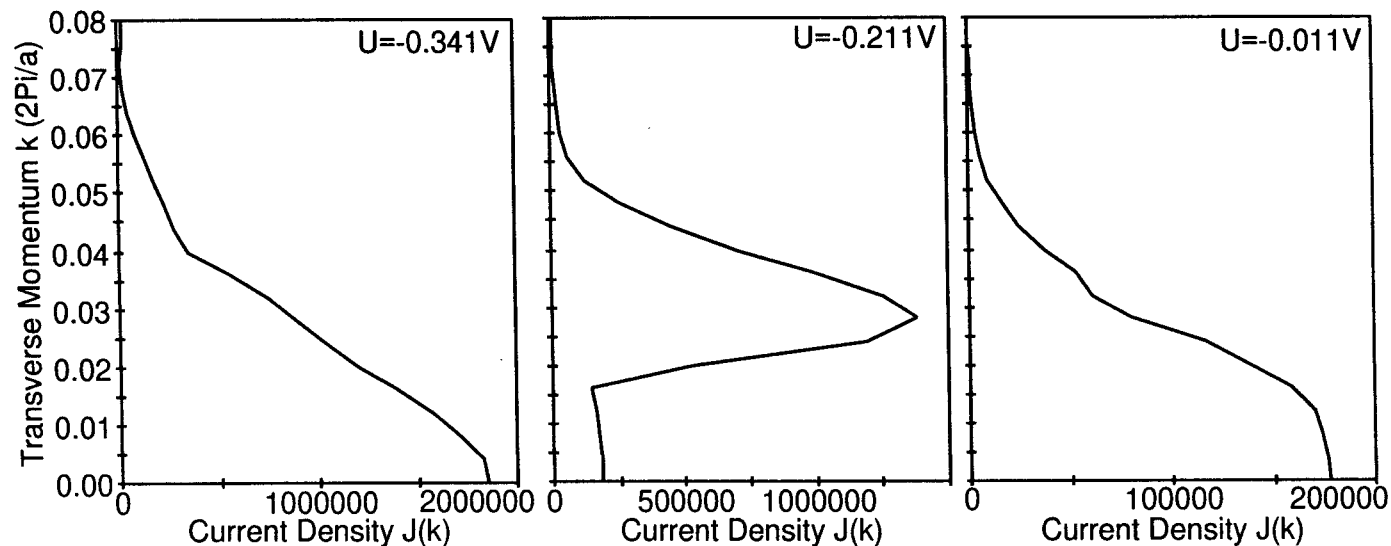
- $J(k)$  decreases with increasing  $k$
- $J(E, k)$  is widely distributed in energy



# Current Voltage Characteristic



- Transverse integration provides qualitatively different results.
- Current dependence of  $k_t$  in  $\langle 100 \rangle$  or  $\langle 110 \rangle$  direction is weak.
- Current may be flowing dominantly outside the zone center.

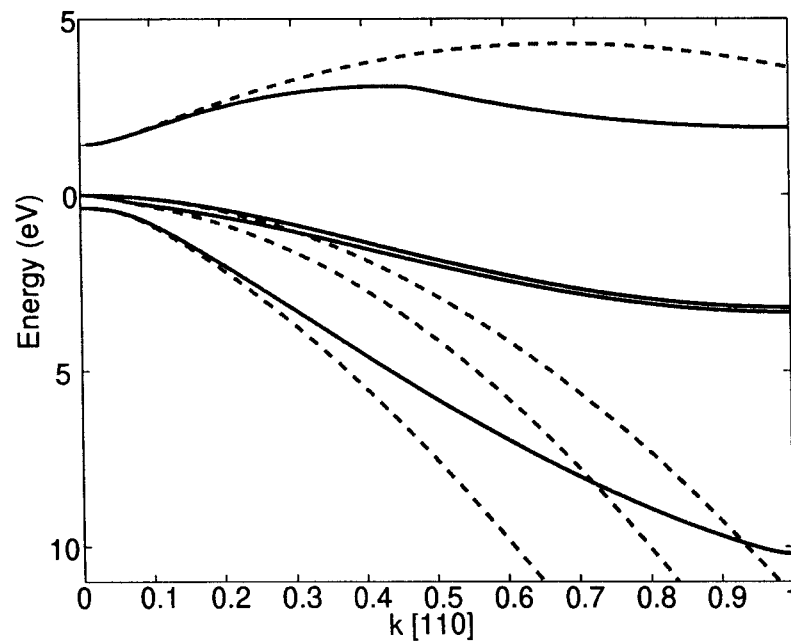


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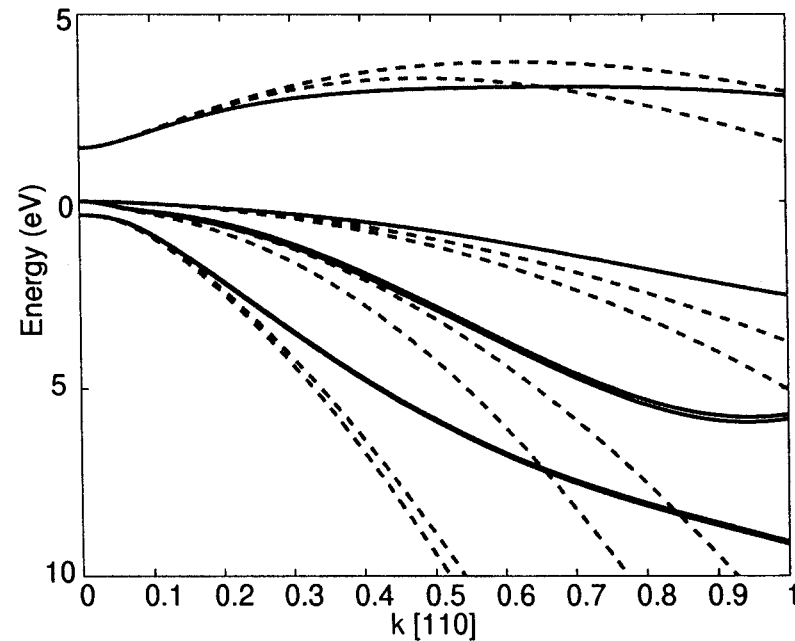
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# Comparison of GaAs $sp^3s^*$ and k.p bulk bandstructure

- 3 Valence bands, 1 conduction band,
- k.p parameters such that bandgaps and effective masses are identical to tight binding



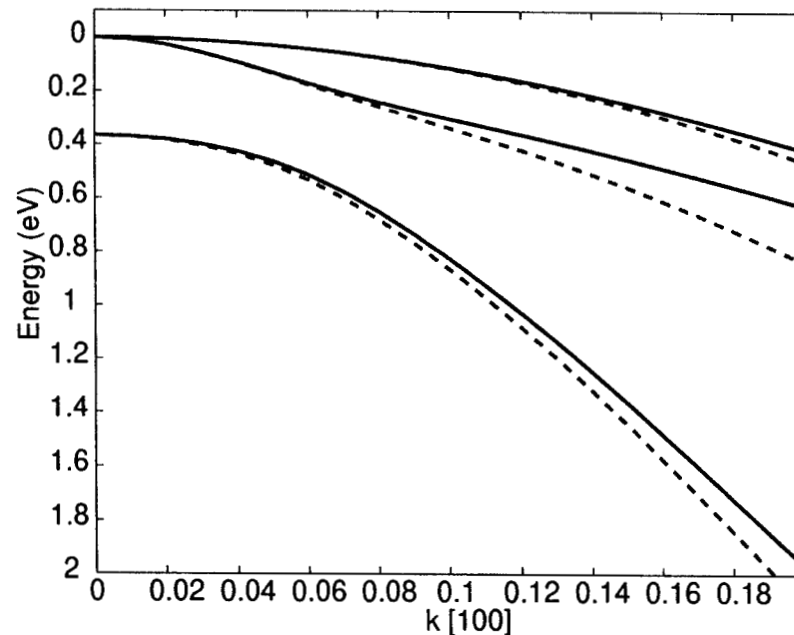
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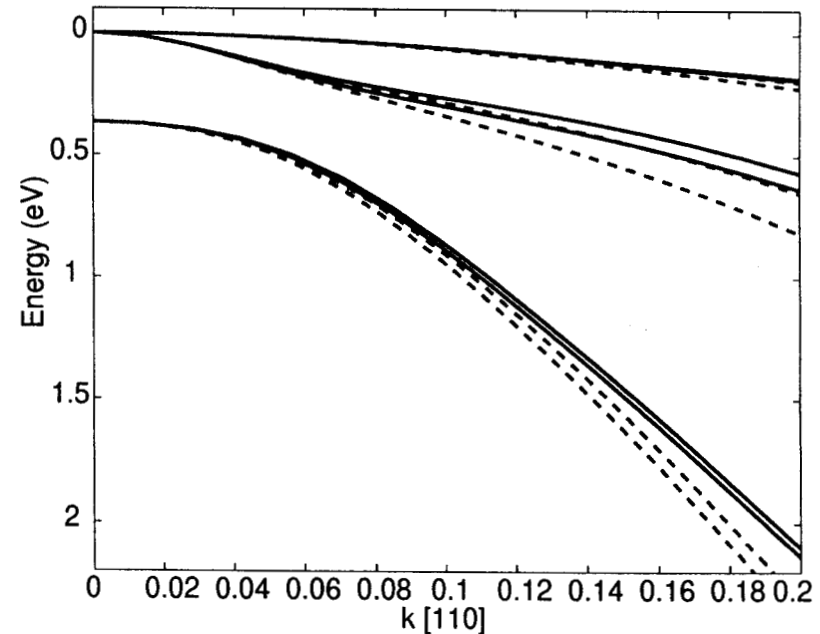
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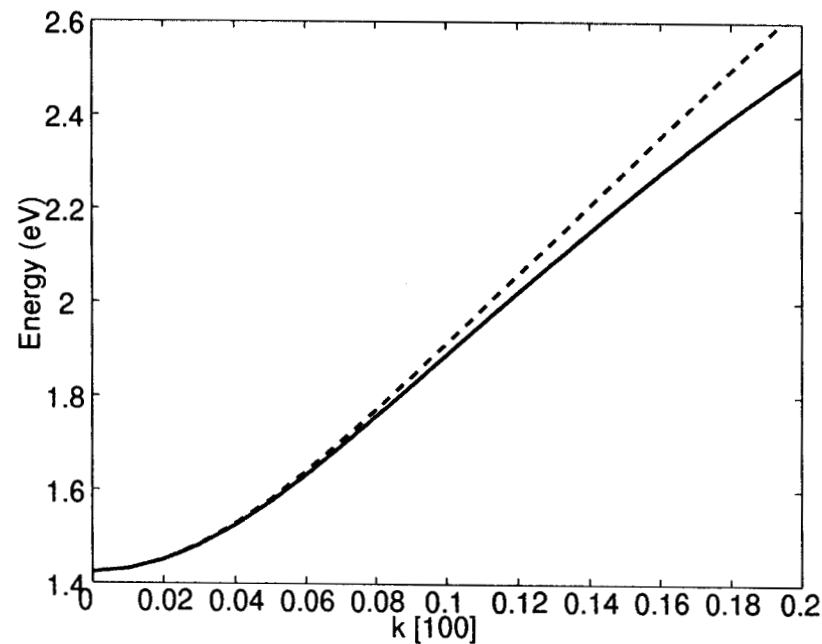
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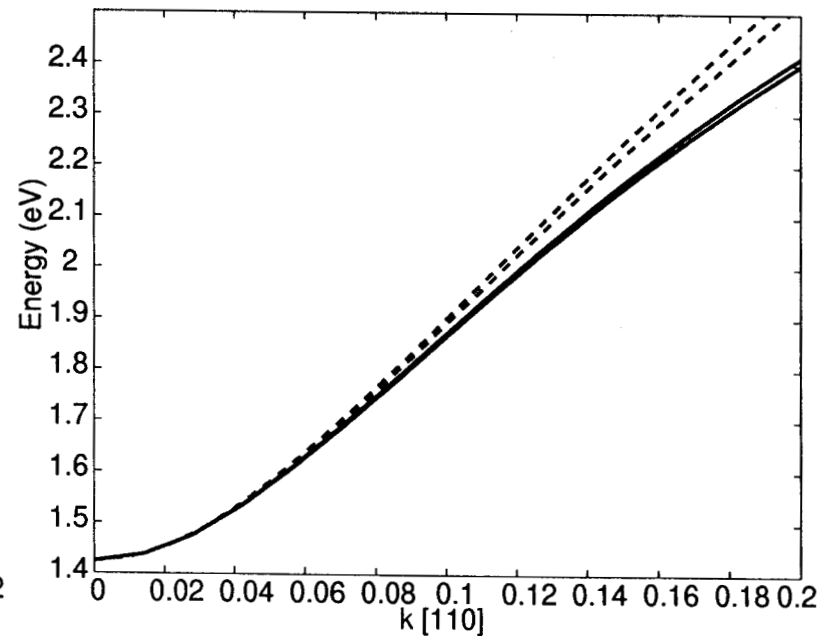
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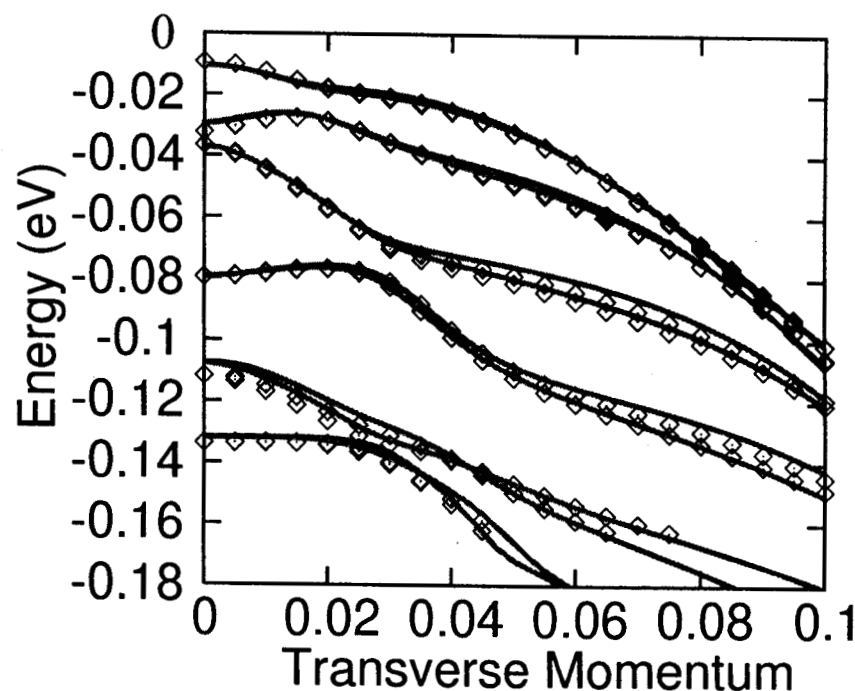
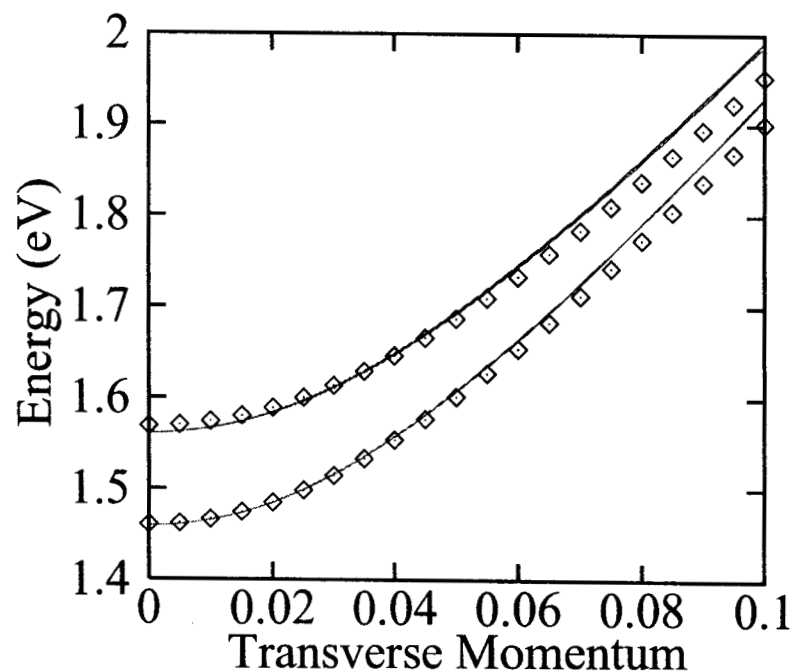


[100]



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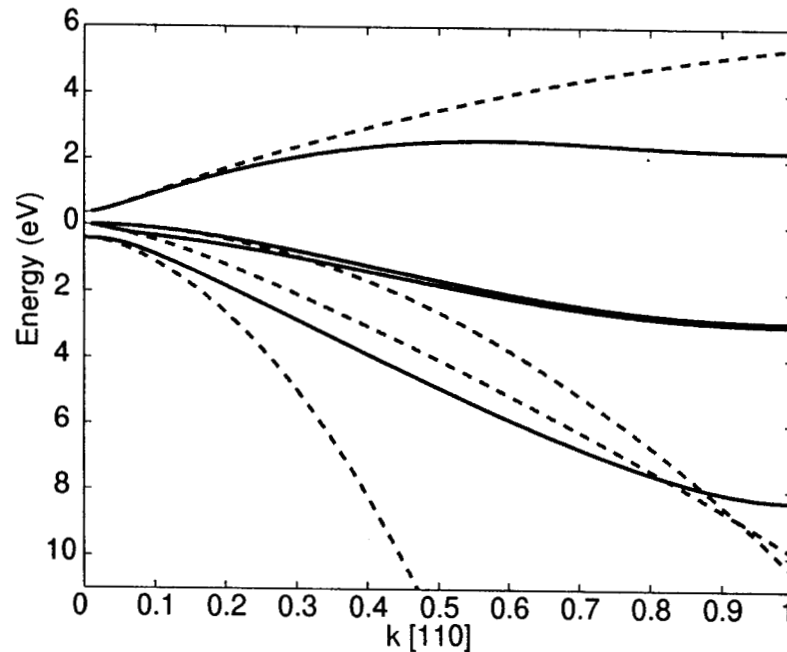
# Sp3s\* vs. k·p: Transverse Subbands in 30ml $\text{Al}_{0.3}\text{Ga}_{0.7}\text{As}/\text{GaAs}/\text{Al}_{0.3}\text{Ga}_{0.7}\text{As}$ Quantum Well



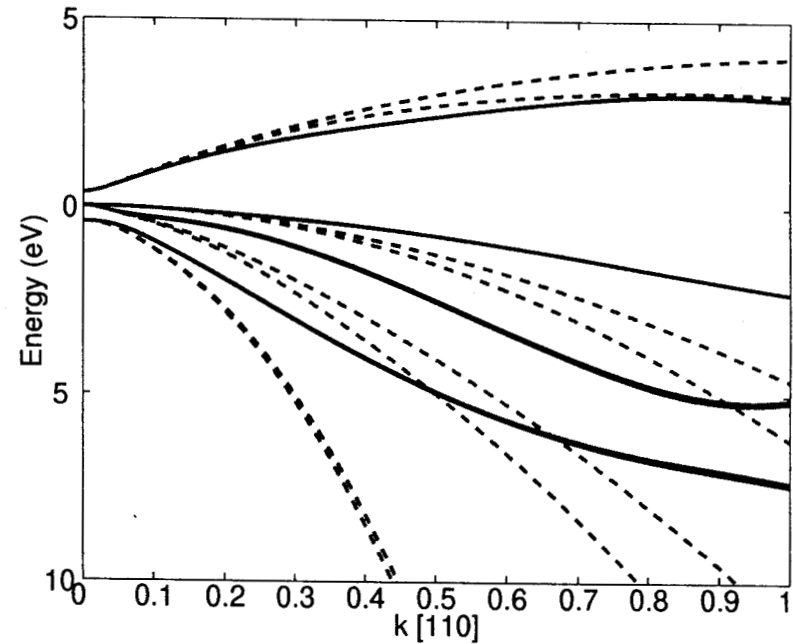
- Subbands in the [100] direction. Lines are k.p, dots are sp3s\*.
- K.p parameters are adjusted to reflect sp3s\* bandgaps and effective masses.

# Comparison of InAs $sp^3s^*$ and k.p bulk bandstructure

- 3 Valence bands, 1 conduction band,
- k.p parameters such that bandgaps and effective masses are identical to tight binding



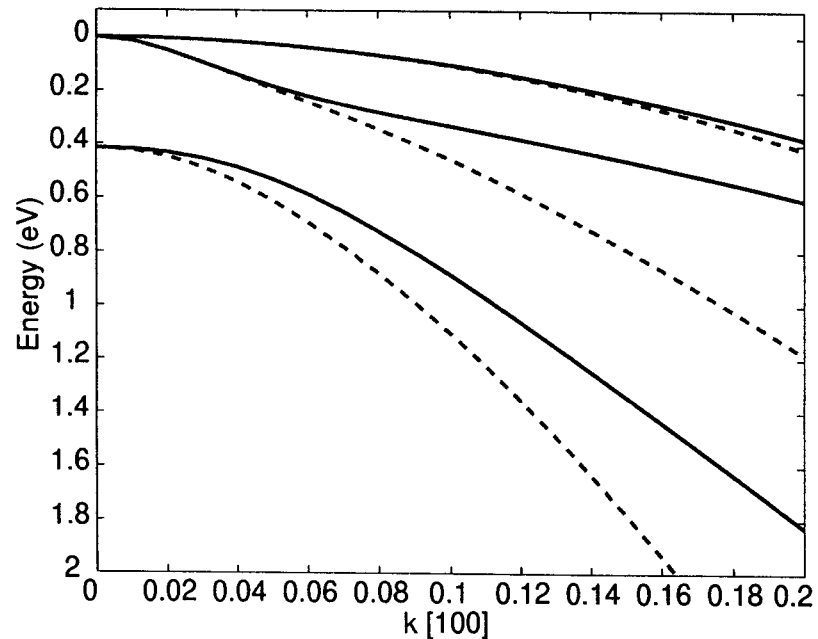
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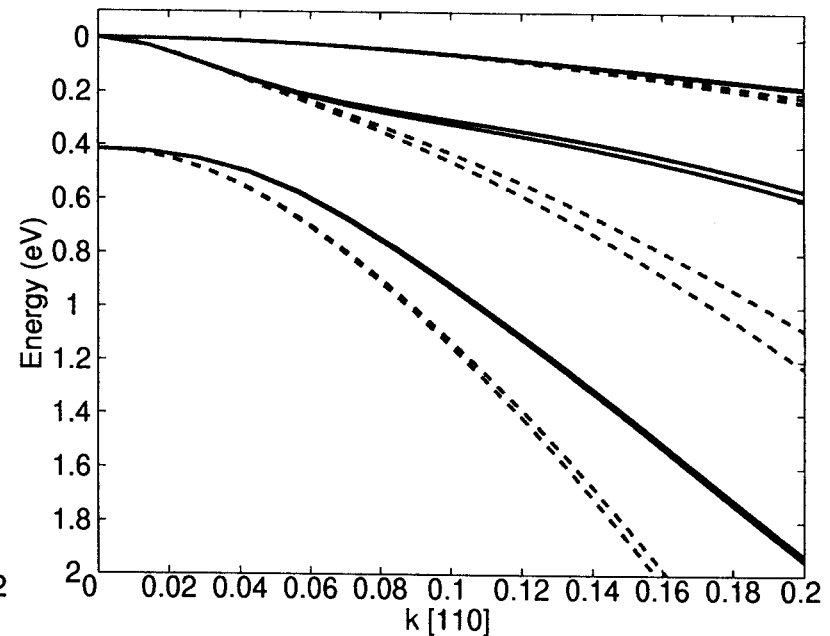
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[100]

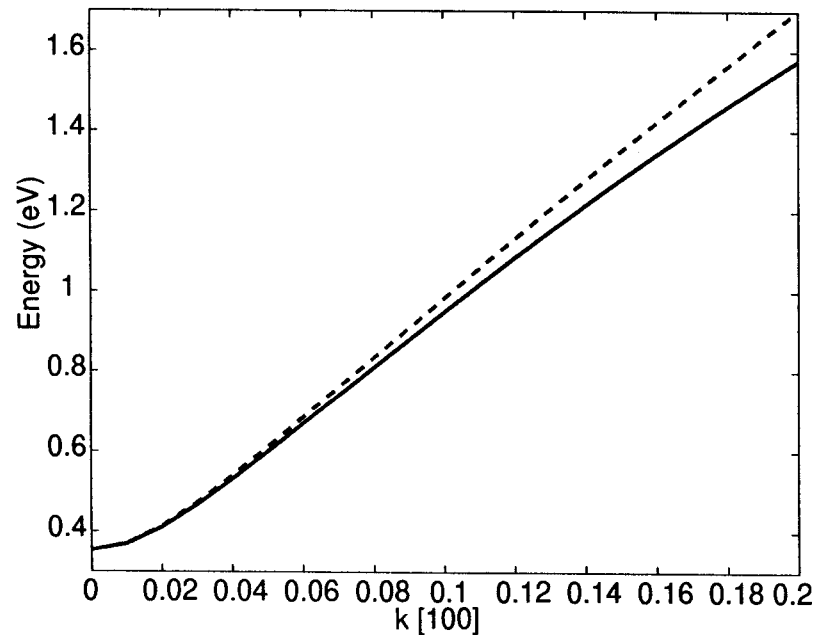


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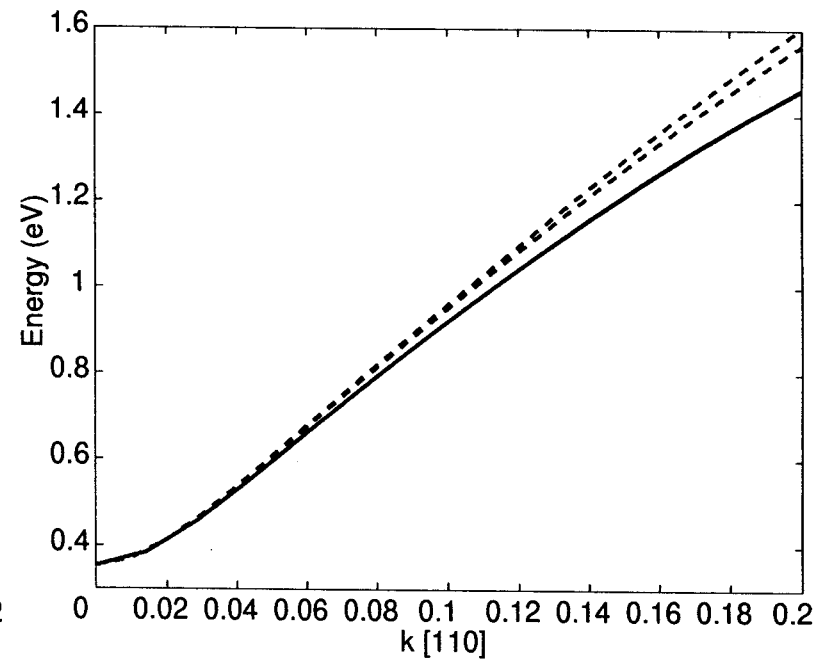


# Comparison of InAs $sp^3s^*$ and k.p bulk bandstructure

- 1 conduction band
- k.p parameters such that bandgaps and effective masses are identical to tight binding



[100]



[110]

## **Open list for pros and cons: tight binding vs. k.p**

- **K.p better than tight binding:**
  - **Easier parameterization.**
  - **Faster simulations.**
- **Tight binding better than k.p:**
  - **Spatial discretization lends itself to incorporation of:**
    - **Space charge effects, arbitrary electrostatic potentials**
    - **Open boundary effects - real structures are typically non-periodic**
    - **Transport simulations (coherent and incoherent)**
  - **Transport away from zone center ( $\Gamma$ -X- $\Gamma$ ), indirect gap materials**
- **Both models fudge:**
  - **TB: adjust interactions to optimize the bands of interest**
  - **K.p: exclude higher bands from basis set and include their effects via Löwdin perturbation.**

# Conclusions

- Full band simulation is essential for:
  - Quantitative electron transport
  - Qualitative hole transport
- We will implement a discretized tight binding model into NEMO and evaluate the performance in transport simulation.
- Genetic algorithm was used to drive NEMO as a black box for structural optimization.